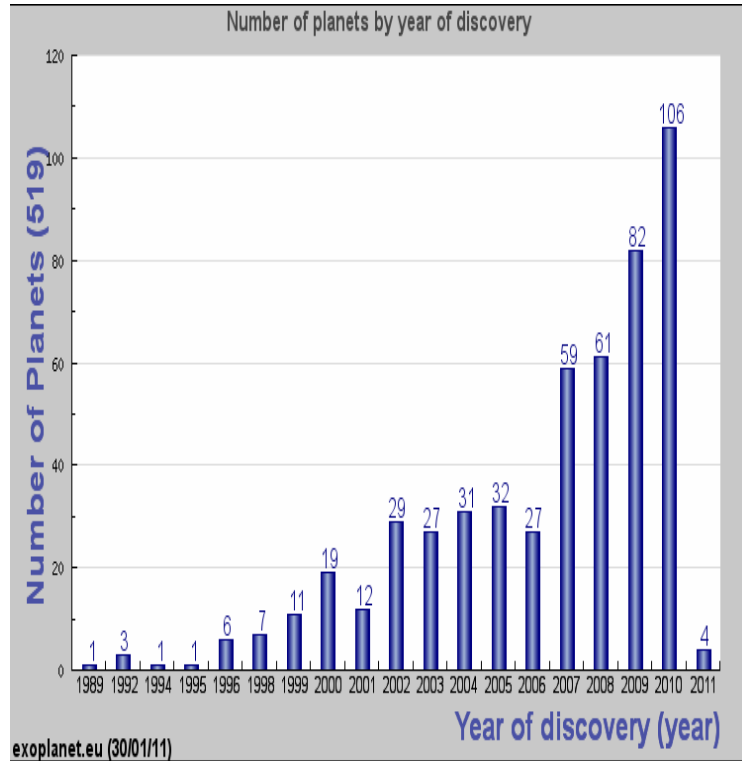


ОБ ОСОБЕННОСТЯХ МАГНИТОСФЕР ГОРЯЧИХ ЭКЗОПЛАНЕТ-ГИГАНТОВ

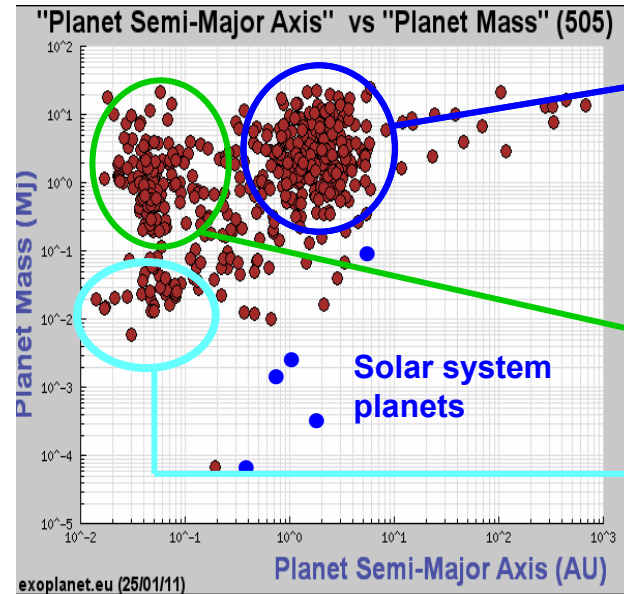
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➤ Exoplanet – Status Jan. 2011



- 435 Exoplanetary systems
- 519 Exoplanets
- 54 Multiple Planetary systems



usual Giants

Hot Jupiters
(0.5-1,75) M_J
136 (27%)

? Evolution of planets
? Formation of terrestrial type worlds

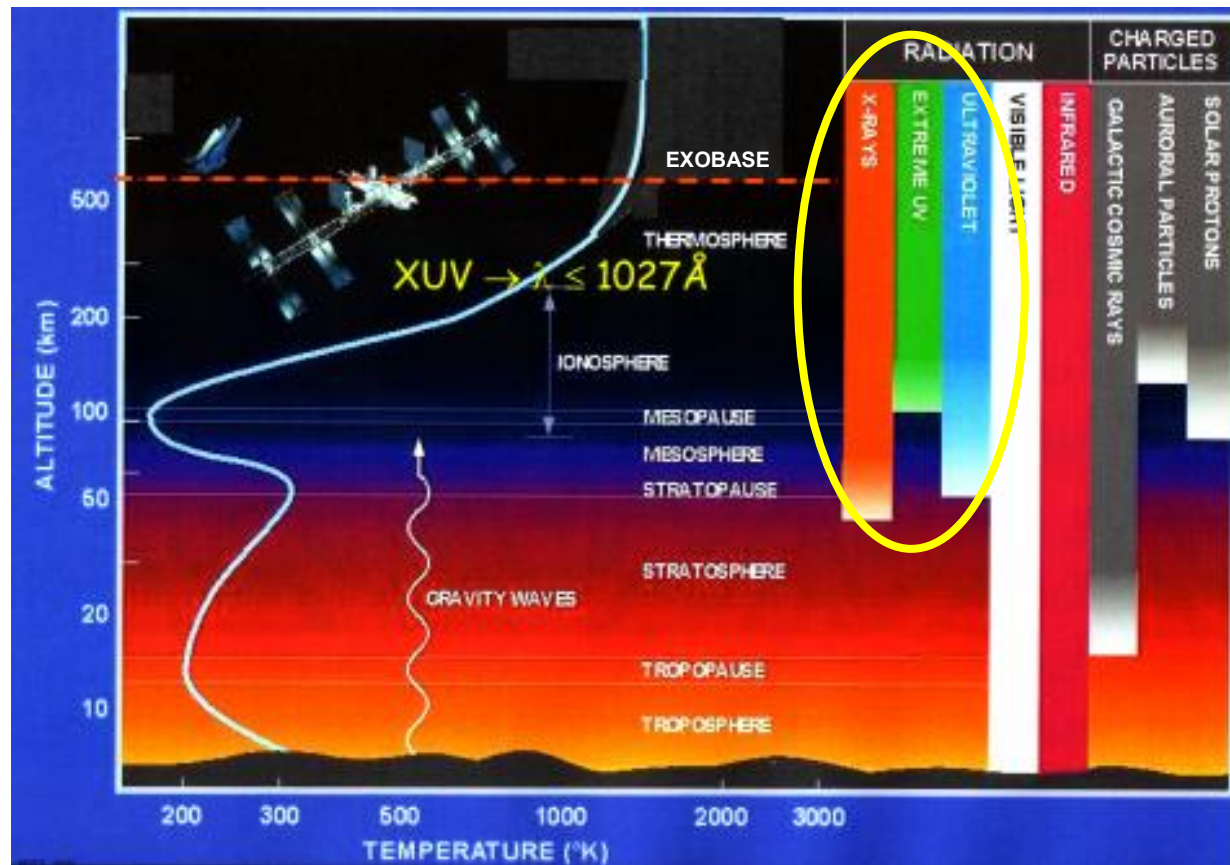
- 28 planets, $< 10 m_{Earth}$
- 9 planets, $< 5 m_{Earth}$



Super-Earths ?

➤ Exoplanet evolution – mass loss

- Stellar X-ray and EUV induced expansion of the upper atmospheres
 - ◆ Stellar XUV luminosity → energy deposition to upper atmospheres



➤ Exoplanet evolution – mass loss

● Soft X-ray and EUV induced expansion of the upper atmospheres

⇒ high *thermal* & *non-thermal* loss rates

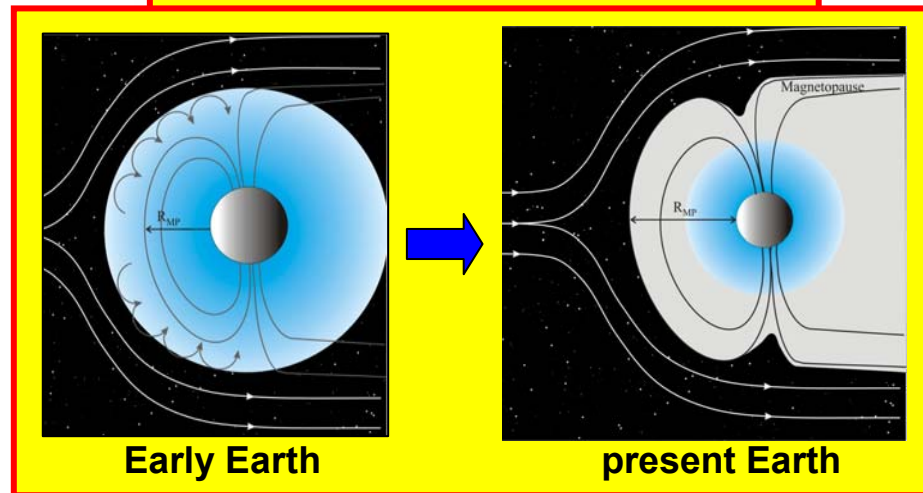
◆ Thermal escape: particle energy $> W_{\text{ESC}}$

→ Jeans escape – particles from “tails”

→ hydrodynamic escape – all particles

The size of magnetosphere is a crucial parameter

Magnetically protected planet



◆ Non-thermal escape:

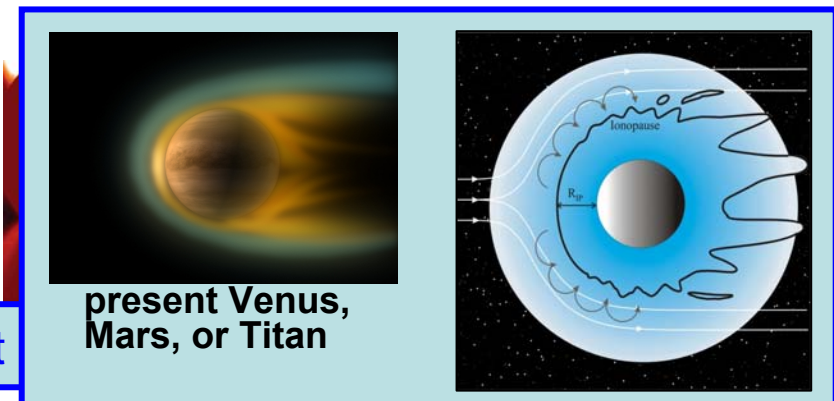
→ Ion pick-up

→ Sputtering (S.W. protons & ions)

→ Photo-chemical energizing & escape

→ Electromagnetic ion acceleration

Magnetically non-protected planet



➤ Exoplanet magnetic fields – role in planet protection

- Magnetic moment estimation from scaling laws → range for possible M

$$M \propto \rho_c^{1/2} \omega r_c^4$$

$$M \propto \rho_c^{1/2} \omega^{1/2} r_c^3 \sigma^{-1/2}$$

$$M \propto \rho_c^{1/2} \omega^{3/4} r_c^{7/2} \sigma^{-1/4}$$

$$M \propto \rho_c^{1/2} \omega^{1/2} r_c^3 \sigma^{-1/2}$$

$$M \propto \rho_c^{1/2} \omega r_c^{7/2}$$



Interval of possible values for
planetary magnetic dipole:

$$M_{max} \dots M_{min}$$

J.-M. Grießmeier, A&A, 2004, 425, 753

J.-M. Grießmeier, Astrobiology, 2005, 5

r_c - radius of the dynamo region (“core radius”): $r_c \sim M_p^{0.75} R_p^{-0.96}$

ρ_c - density in the dynamo region

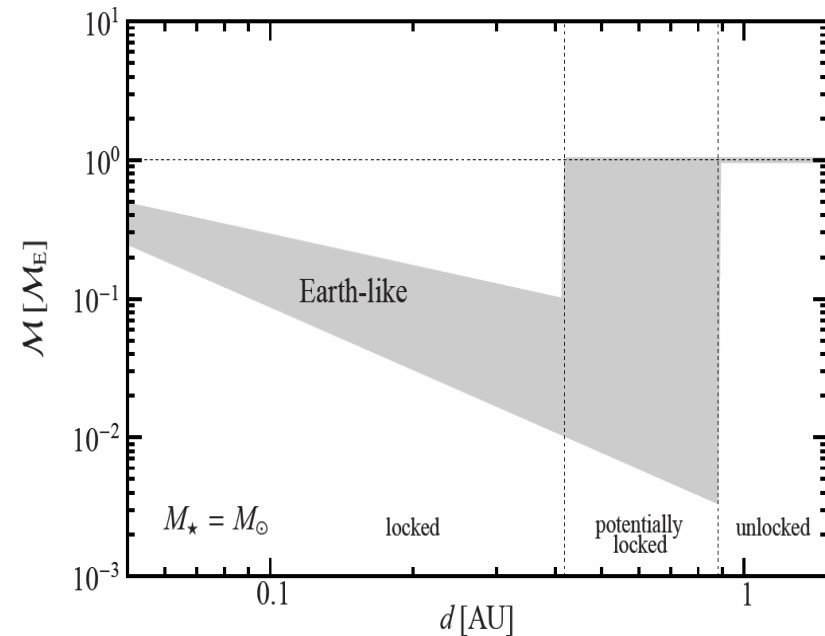
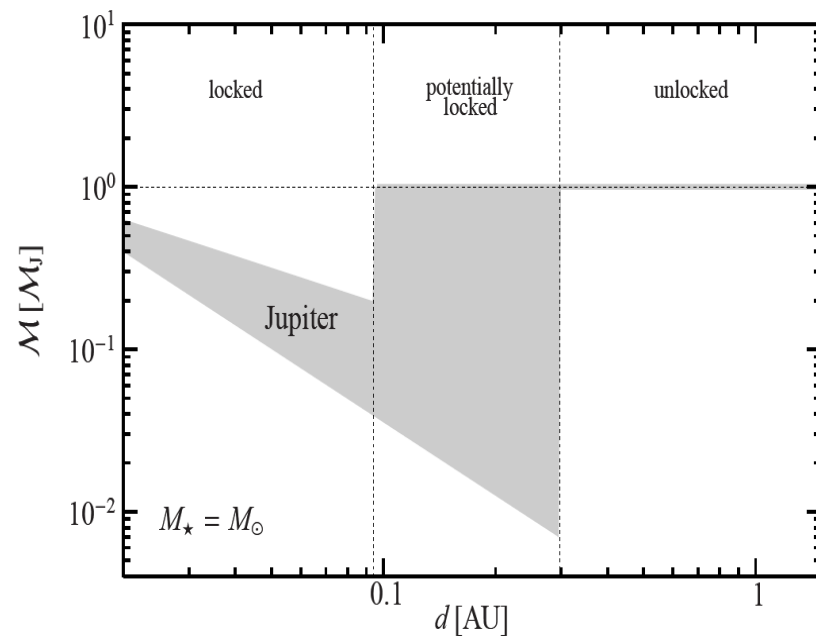
σ - conductivity in the dynamo region

ω - planet angular rotation velocity

➤ Exoplanet magnetic fields – role in planet protection

- Magnetic moment estimation from scaling laws → range for possible M

- ◆ Limitation of M by tidal locking [Grißmeier, J.-M., et al., *Astrobiology*, 5(5), 587, 2005]



Tidal locking ⇒ strongly reduced magnetic moments

➤ Exoplanet magnetic fields – role in planet protection

● Magnetospheric protection of evaporating/eroded planetary atmospheres

Khodachenko et al., PSS, 55, 631, 2007; Khodachenko et al., Astrobiology, 7, 167, 2007

◆ CME induced H⁺ ion pick-up loss at 0.05 AU for ‘Hot Jupiters’ → [HD209458 b](#)

Conditions	S [s ⁻¹]	L [g s ⁻¹]	\mathcal{M} [\mathcal{M}_{Jup}]	n_{CME} [cm ⁻³]	r_s [r_{pl}]	Γ [M_{pl}]
CME _{min} , \mathcal{M}_{max}	9×10^{34}	<u>1.5×10^{11}</u>	0.1	6300.0	2.33	1.56×10^{-2}
CME _{max} , \mathcal{M}_{max}	7×10^{37}	2×10^{13}	0.1	7.5×10^4	1.54	0.2
CME _{min}	7.2×10^{36}	1.2×10^{13}	0.017	6300.0	1.3	0.12
CME _{max}	8.2×10^{37}	1.37×10^{14}	0.059	7.5×10^4	1.3	1.43
CME _{min}	8.4×10^{37}	1.4×10^{14}	0.012	6300.0	1.15	1.46
CME _{max}	9.5×10^{38}	1.6×10^{15}	0.041	7.5×10^4	1.15	17.0
CME _{min} , \mathcal{M}_{min}	5.0×10^{39}	8.35×10^{15}	0.005	6300.0	1.0 [0.85]	87.0
CME _{max} , \mathcal{M}_{min}	5.7×10^{40}	9.5×10^{16}	0.005	7.5×10^4	1.0 [0.56]	990.0

Mass loss $\sim 10^{11}$ g/s even for weak CMEs & $\mathcal{M}_{\text{max}} \Rightarrow$ strong magn. protection in reality

➤ Exoplanet magnetospheres – importance of magnetodisk

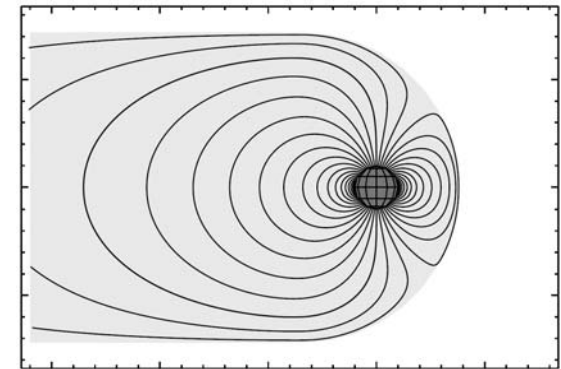
- ◆ **Relatively large amount of observed Hot Jupiters (27%):**
"survival" of close-in giants indicates their efficient protection against of extreme plasma and radiation conditions

- ◆ **All estimations were based on too simplified model**

Magnetospheric protection of exoplanets was studied assuming a simple planetary dipole dominated magnetosphere

→ dipole mag. field $\mathbf{B} = \mathbf{B}_{\text{dip}} \sim M / r^3$ balances stellar wind ram pressure

→ big M are needed for the efficient protection (but tidal locking → small M → small R_s)



J.-M. Grießmeier, A&A, 2004, 425, 753

- ◆ **Specifics of close-in exoplanets → [new model](#)**

→ strong mass loss of a planet should lead to *formation of a plasma disk* (similar to Jupiter and Saturn) → Magnetodisk dominated magnetosphere

→ more complete planetary magnetosphere model, including the whole complex of the magnetospheric electric current systems

➤ Exoplanet magnetospheres – importance of magnetodisk

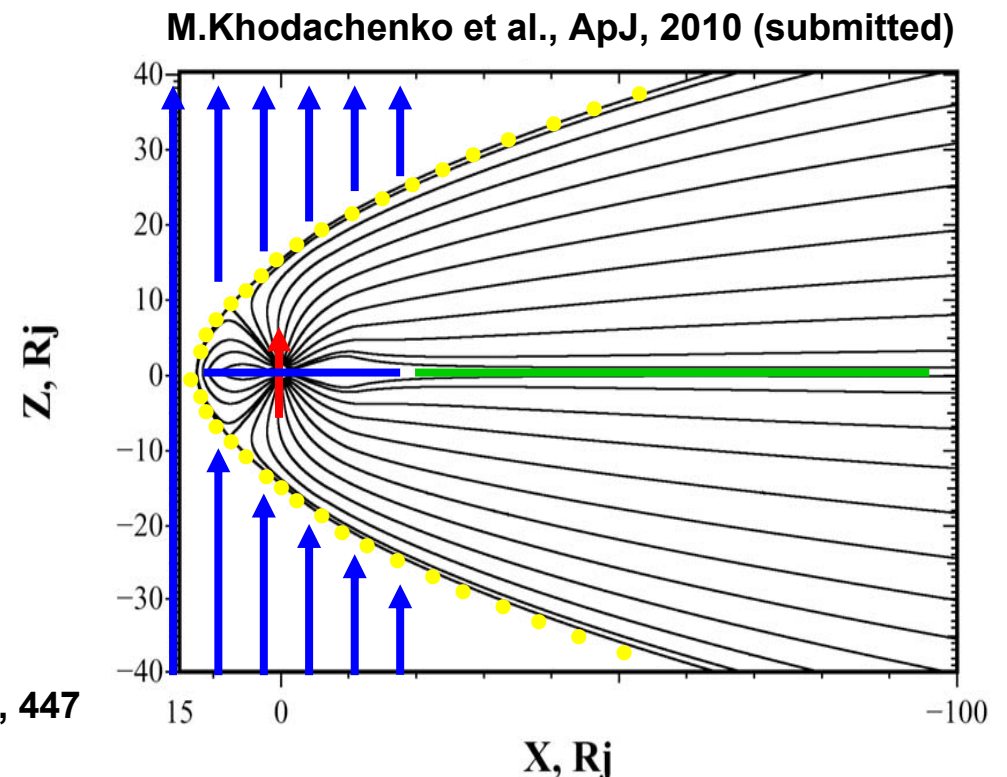
● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

Semi-analytical model. Key assumption: magnetopause is *approximated by paraboloid of revolution* along planet-star (V_{SW}) line

◆ PMM considers mag. field of different current systems on the boundaries and within the boundaries of a planetary magnetopause:

- planetary *magnetic dipole*;
- current system of *magnetotail*;
- *magnetodisk*;
- *magnetopause* currents;
- *magnetic field of stellar wind*, partially penetrated into the magnetospheric obstacle.

I.Alexeev, Geomag.&Aeronomia, 1978, 18, 447



➤ Exoplanet magnetospheres – importance of magnetodisk

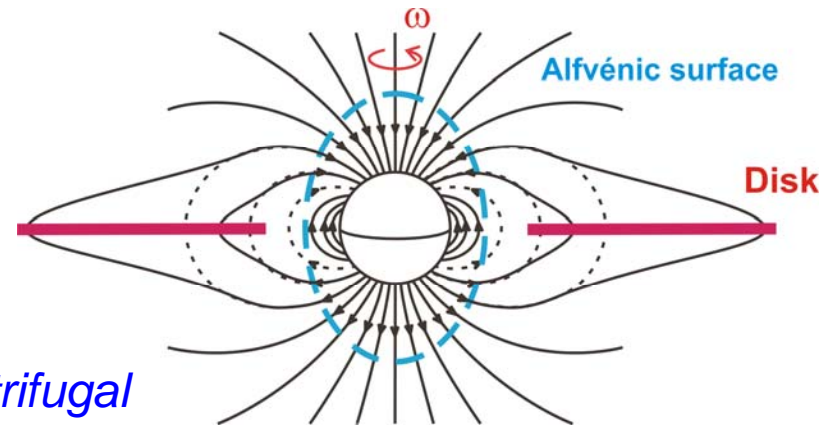
● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

◆ Formation of magnetodisk

→ “sling” model:

dipole mag. field can drive plasma
in co-rotation regime only inside

“**Alfvénic surface**” ($r < R_A$); → **Centrifugal**
escape of plasma

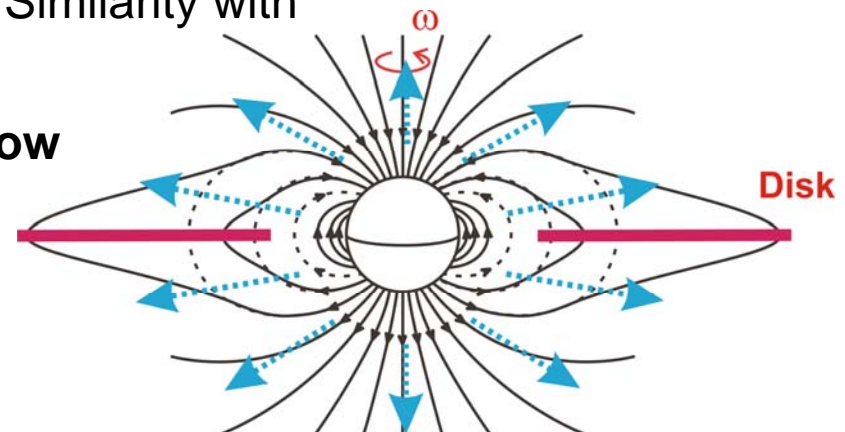


→ “material-escape driven” models: → **Hydrodynamic** escape of plasma

(a) **Fully ionized plasma outflow** – Similarity with
heliospheric current sheet (disk)

(b) **Partially ionized material outflow**

Background magn.field (dipole),
charge separation electric field,
ambipolar diffusion, azimuth.Hall
current in equator.plane



➤ Exoplanet magnetospheres – importance of magnetodisk

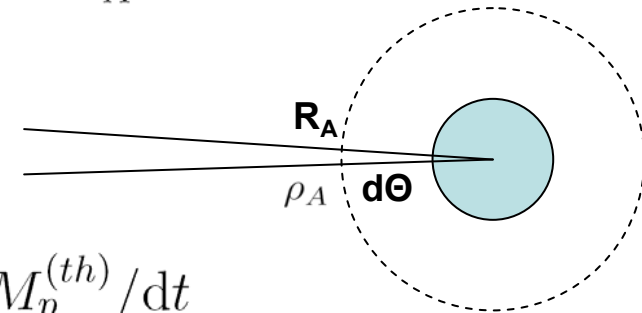
● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

◆ Formation of magnetodisk – “sling” model (analogy with the Jupiter):

→ “Alfvénic surface” radius R_A

Co-rotation until $\varepsilon_p = \frac{\rho_A \omega_p^2 R_A^2}{2} = \varepsilon_B = \frac{M_d^2}{2\mu_0 R_A^6}$, $M_d = \mu_0 \mathcal{M} / (4\pi)$

⇒ R_A :
$$R_A = \left(\frac{M_d^2}{\mu_0 \rho_A \omega_p^2} \right)^{1/8}$$



→ ρ_A is estimated from thermal mass loss $dM_p^{(th)} / dt$

$$\rho_A = \frac{dM_p^{(th)} / dt}{2\pi R_A^3 \omega_p \delta\theta} \Rightarrow \frac{R_A}{r_p} = \left(\frac{2\pi \delta\theta B_{d0}^2 r_p}{\mu_0 \omega_p (dM_p^{(th)} / dt)} \right)^{1/5}$$

Increase of ω_p and $dM_p^{(th)} / dt$ → decrease of R_A

➤ Exoplanet magnetospheres – importance of magnetodisk

● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

◆ Magnetic field structure in PMM with magnetodisk

→ $r < R_A$: magnetic field of dipole ($\sim R^{-3}$)

$$\mathbf{B}(r, z = 0) = B_\theta(r) \theta_0 = B_{d0} r_p^3 / r^3 \theta_0$$

→ $r > R_A$: conservation of m.flux reconnected across the disc
 ⇒ m.field of the disk (and current density) $\sim R^{-2}$

$$\mathbf{B}(r, z \sim 0) = B_r(r) \mathbf{r}_0 = B_{d0} r_p^3 / (R_A r^2) \mathbf{r}_0$$

◆ Determination of sub-stellar magnetopause distance:

pressure balance condition

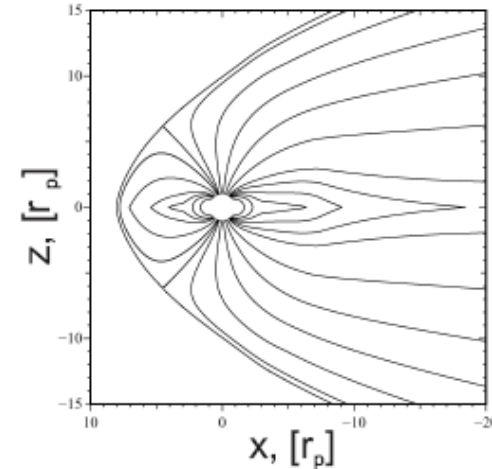
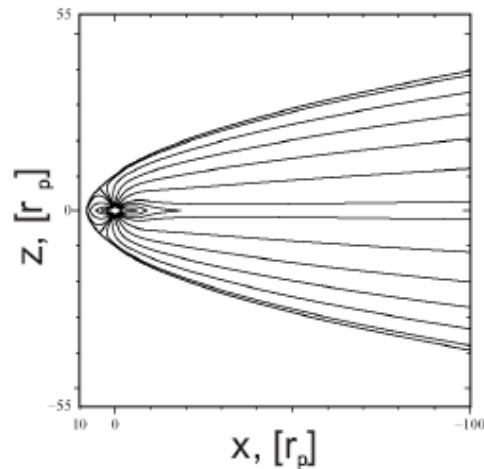
$$\rho_{sw} \tilde{v}_{sw}^2 = \frac{\kappa^2 (B_{ds} + B_{MDS})^2}{2\mu_0} + p_{mp} \Rightarrow R_s \sim \frac{B_{d0}^{1/2} r_p^{3/2}}{(2\mu_0 p_{sw})^{1/4} R_A^{1/2}} (1 + \kappa^2)^{1/4}$$

→ Decrease of R_A ⇒ increase of R_s

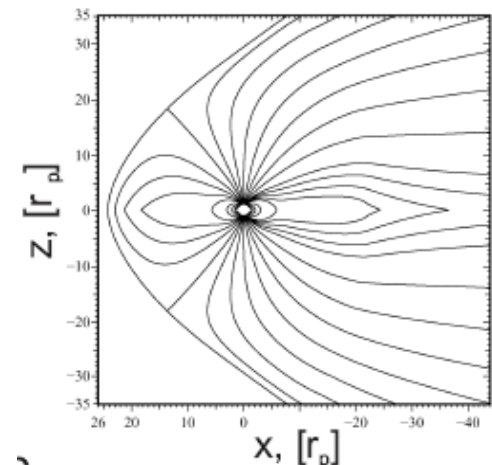
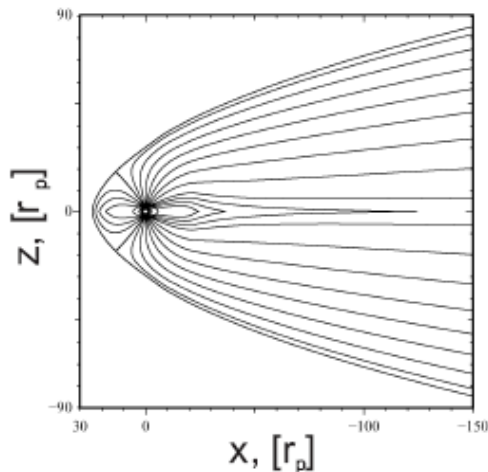
➤ Exoplanet magnetospheres – importance of magnetodisk

● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

◆ Magnetosphere at **0.045 AU**, $R_S = 8.0 R_J$ (tidally locked)



◆ Magnetosphere at **0.3 AU**, $R_S = 24.2 R_J$ (tidally un-locked)



➤ Exoplanet magnetospheres – importance of magnetodisk

● Paraboloid Magnetospheric Model (PMM) for ‘Hot Jupiters’

◆ magnetospheric parameters (estimated and calculated)

Table 4: “Hot Jupiter” magnetopause stand-off distance at substellar point R_s and its major control parameters. $R_s^{(dip+MD+tail)}$ is given by PMM taking into account screened planetary dipole, screened magnetodisk and magnetotail currents; $\tilde{R}_s^{(dip+MD)}$ is non-self-consistent estimate of R_s followed from (18); $R_s^{(dip)}$ is an estimate of R_s for the case of only a dipole magnetosphere given by (3). ¹: Tidally locked. ²: Not tidally locked. ³: Jupiter

d	$R_s^{(dip+MD+tail)}$	$\tilde{R}_s^{(dip+MD)}$	$R_s^{(dip)}$	R_A	$\frac{\mathcal{M}_{MD}}{\mathcal{M}}$	ω_p	$\frac{dM_p^{(th)}}{dt}$
[AU]	$[r_p]$	$[r_p]$	$[r_p]$	$[r_p]$		$[\omega_J]$	$[g\ s^{-1}]$
0.045 ¹	8.0	9.27	5.76	3.1	1.64	0.118	1.06×10^{10}
0.1 ¹	8.27	9.06	6.16	4.42	1.29	0.036	1.80×10^9
0.3 ²	24.2	25.6	15.0	6.97	3.59	1.0	1.84×10^8
5.2 ³ (Jupiter)	71.9	69.3	41.8	19.8	3.32	1.0	1.0×10^6

➤ Conclusions

Magnetodisks of close-in giant exoplanets (Hot Jupiters) influence the structure and character of their magnetospheres, leading to a new type of „*magnetodisk dominated*“ magnetosphere.

Extended up to **(40 – 70) %** magnetodisk magnetospheres, as compared to dipole type ones, may efficiently protect planetary environments, even close to a host star.